

# Nuclear data uncertainty propagation

# with Monte Carlo methods

D. Rochman,

# . Koning, S.C. van der Marck and D.F. daCruz

Nuclear Research and Consultancy Group,

NRG, Petten, The Netherlands

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All slides can be found at:

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ftp://ftp.nrg.eu/pub/www/talys/bib\_rochman/presentation.html.

### Who are we ?





- Solution NRG: a leading nuclear sector service provider
- Over 50 years experience in nuclear technology
- Over 400 employees (10 in R&D reactor physics and simulations)
- Turnover approximately Meuros 60 / per year
  - High Flux Reactor, Hot Cell Laboratories and Radiological labs





### What are nuclear data?

The term "nuclear data" can have different meaning,

- dusty books, constants, mature field, code inputs,
- list, Schrodinger equation, unexciting...
- but this is not ! (I'm going to prove that)









## Why are nuclear data important ? (Part 1)

Better nuclear data can help for:

- ► safety margins, fuel storage,
- ► life-time extension,
- cost reduction in design of new systems,
- ► isotope production,
- ► safety of people (shielding),
- ► waste transmutation,
- development of future systems.

Better nuclear data have a limited effect on:

- current reactor operation,
- current reactor safety,...
- ◄ accident simulation,
- ◄ proliferation,
- ◀ Chernobyl, TMI, Fukushima.

#### Leistungszuwachs seit 1955

Elektrische Nettoleistung der Kernkraftwerke weltweit von 1955 bis 2011 in Megawatt (MW)

#### Augmentation de la puissance depuis 1955

Puissance électrique nette des centrales nucléaires dans le monde de 1955 à 2011 en mégawatts (MW)



### Why are nuclear data important ? (Part 2)



### Nuclear data uncertainties: general comments

- uncertainties are not errors (and vice versa),
- l they are related to risks, quality of work, money, perception, fear, safety...

Uncertainty  $\leftrightarrows$  safety  $\rightleftharpoons$  professionalism

- III True uncertainties do not exist ! They are the reflection of our knowledge and methods.
- $\blacksquare$  All the above for covariances
- **#**| The importance of nuclear data uncertainties should be checked. If believed

Our motivation: Any justification for not providing uncertainties should become obsolete

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## Mission and objective of the NRG Physics Team

# Our mission: improve nuclear simulations





"Towards sustainable nuclear energy: Putting nuclear physics to work",

A.J. Koning and D. Rochman, ANE 35 (2008) 2024.

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## TMC for nuclear data uncertainty propagation, what else ?

- $\bigcirc$ + No covariance matrices (no 2 Gb files) but every possible cross correlation included,
- $\bigcirc$ + No approximation but true probability distribution,
  - + Only essential info for an evaluation is stored,
  - + No perturbation code necessary, only "essential" codes,
  - + Feedback to model parameters,
  - + Full reactor core calculation and transient,
- $\bigcirc$ + Also applicable to fission yields, thermal scattering, pseudo-fission products, all isotopes (...just everything),
- $(\dot{})$ + Other variants: AREVA (NUDUNA), GRS (XSUSA), CIEMAT (ACAB), PSI (NUSS), CNRS Grenoble..., based on covariance files,  $\bigcirc$ 
  - + Many spin-offs (TENDL covariances, sensitivity, adjustment...)
  - + Computer time (not human time)
  - + QA.
  - Needs discipline to reproduce,
  - Memory and computer time (not human time),
  - Need mentality change.

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Angle (deg)

Energy (MeV)



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### **Considered data in TMC (or fast TMC)**

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Several hundreds of random ENDF files for transport + depletion

- 3 Major actinides: <sup>235</sup>U, <sup>238</sup>U, <sup>239</sup>Pu,
- Light elements: lighter than oxygen,
- Thermal scattering data: H in H<sub>2</sub>O, D in D<sub>2</sub>O, C in Carbon,
- All Fission yields (e.g. <sup>234,235,236,238</sup>U, <sup>239,240,241</sup>Pu, <sup>237</sup>Np, <sup>241,243</sup>Am, <sup>243,244</sup>Cm),
- All Minor actinides (e.g. <sup>234,236,237</sup>U, <sup>237</sup>Np, <sup>238,240,241,242</sup>Pu, Am, Cm),
- All fission products (e.g. from Ge to Er), and decay data,

(fast) TMC can be applied to any input data, propagating uncertainties to any outputs



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(fast) TMC can be applied to any input data, propagating uncertainties to any outputs

TMC was already applied to

- criticality-safety, shielding, pincell/assembly burn-up, full core, activation,
- PWR, BWR, Gen-IV systems,
- UO<sub>2</sub>, MOX fuels,
  - MCNP, SERPENT, FISPACT, DRAGON, PANTHER, RELAP-5
# Application: thermal scattering for H in $H_2O$ or $S(\alpha,\beta)$ tables (with MCNP)

Random parameters of the  $S(\alpha,\beta)$  for inelastic scattering



## Systematical study on UO<sub>2</sub>/MOX assembly uncertainties



- Different UO<sub>2</sub>/MOX assemblies (PWR, BWR, VVER, AGR, CANDU, fast systems),
- Burn-up calculated with SERPENT,
- All major nuclear data taken into account.
- $\implies$  systematical study on  $k_{eff}$ , inventory, heat...

#### **Comparison of** $\Delta \mathbf{k}_{\infty}$ **for assemblies and full core (SERPENT)**



# TMC applied to PWR assembly burn-up calculations with DRAGON N



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#### TMC applied for burn-up calculations: decay heat



## TMC applied for burn-up calculations: decay heat uncertainties



# Effect of H in H<sub>2</sub>O for a full core PWR (courtesy of O. Cabellos, UPM, Spain)

Method: TMC applied to COBAYA (3D multigroup core calculations) + SIMULA (coupled neutronic-thermohydraulics 3D core calculations)

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UAM7 – Paris (France), April 10-12, 2013

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If we can do a calculation once, we can also do it a 1000 times, each time with a varying data library.

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Well, then uncertainty propagation with TMC takes  $1000 \times 1000$  longer than a single calculation...

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There is a solution with Monte Carlo codes: (in fact 2 solutions)

- ✤ fast GRS method,
- ✤ and fast TMC.

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"Efficient use of Monte Carlo: uncertainty propagation",

D. Rochman, W. Zwermann et al., submitted to NSE, May 2013.



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# 2012: fast GRS method

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First presented in PHYSOR-2012 by W. Zwermann *et al*.. It takes advantage of conditional expectations:

If two output variables  $k^{(1)}$  and  $k^{(2)}$  are identically distributed and conditionally independent given the vector of nuclear data input then

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#### In practice:

- 1. perform i = 1..500 MCNP short calculations with random nuclear data and a fixed seed  $s_1 \Longrightarrow k_{eff}^{(1)}(i)$
- 2. repeat for j = 1..500, same random nuclear data but fixed seed  $s_2 \implies k_{eff}^{(2)}(j)$

There is no necessity to have small  $\sigma_{statistics}$  !! each run can be (very) short

#### fast GRS method

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 $2 \times 500$  "short" runs ~  $2 \times$  "long" run in time

If a single calculation takes *m* histories ( $\sigma_{stat}$  small enough), then repeat it *n* times with *m/n* histories, random nuclear data and random seeds.

 $\sigma_{\text{total}}^2 = \sigma_{\text{statistics}}^2 + \sigma_{\text{nuclear data}}^2$  still holds.



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•		• •			•	
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	• •		• •			•
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	n runs		$\left\{ \begin{array}{c} \sigma(\overline{k}) \\ \sigma_{total}^2 \end{array} \right.$	$\sim \sigma_{\text{stat}}$ = $\frac{1}{n-1} \sum_{i=1}^{n} $	T sec. $\int_{1}^{1} \left( k_{i} - \overline{k} \right)^{2} \checkmark$	
GF	S 2013					D. Rochman – 25 / 33

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• •		• • •			• • •		
run n	nuclear data <i>n</i>	seed $s_n$	m/n hist.	T/n sec.	$\mathbf{k}_n \pm \mathbf{\sigma}_n$		
n runs		$\begin{cases} \sigma^2_{total} \\ \sigma^2_{statistic} \end{cases}$	$= \frac{1}{n-1} \sum_{i=1}^{n} \sum_{i=1}^{n} \sigma_{i}$	T sec. $\int_{1}^{2} (k_i - \overline{k})^2 \checkmark$			
					D. Rochman – 25 / 33		

## The fast methods

- $\odot$  as fast as S/U methods (1-2 × longer than 1 single calculation),
- $\odot$  tested on criticality & shielding benchmarks, burn-up (k<sub>eff</sub> and inventory),

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- $\odot$  as fast as S/U methods (1-2 × longer than 1 single calculation),
- $\odot$  tested on criticality & shielding benchmarks, burn-up (k<sub>eff</sub> and inventory),
- © Example: the Martin-Hoogenboom benchmark

MCNP6 model: 241 fuel assemblies,



- \*  $357 \times 357 \times 100$  regions  $(1.26 \times 1.26 \times 3.66 \text{ cm}^3)$ : 6.4 million cells for generated power (f7)
- \* 1 calculation takes  $2 \times 10^{11}$  histories ( $\sigma_{statistics} = 0.25$  % at the center, 500 weeks on 1 cpu)

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- \* 1 calculation takes  $2 \times 10^{11}$  histories ( $\sigma_{\text{statistics}} = 0.25$  % at the center, 500 weeks on 1 cpu)
- ➤ Uncertainty on local pin power due to <sup>235</sup>U, <sup>238</sup>U, <sup>239</sup>Pu and H in H<sub>2</sub>O thermal scattering in each cell ?



- \* TMC: 500 random runs of  $2 \times 10^{11}$  histories (500 weeks for each),
- ✤ TMC is not applicable,
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### (1) Source convergence

- ➡ For both methods, a first calculation is run with fixed nuclear data to obtain a reasonably converged fission source.
  - All subsequent short simulations start with this fission source:
    - each with 10 inactive cycles and 90 active cycles of  $4 \times 10^6$  histories, and random nuclear data,
    - source convergence tested with the MCNP6 built-in indicator (fission source entropy).
    - $\checkmark$  362 short runs out of 508 were then accepted for fast TMC.
    - $\sim$  2 × 122 short runs out of 2 × 328 were then accepted for fast GRS.



# (2) Statistical uncertainty estimation

- \* In MCNP eigenvalue calculation,  $\sigma_{\text{stat}}$  is usually underestimated.
- \* An independent estimation of  $\sigma_{stat}$  is therefore necessary for fast TMC,
- ✤ From the 508 short runs, the first 389 were repeated with fixed nuclear data,
- \* 274 were then accepted due to source convergence.



# (2) Statistical uncertainty estimation

- \* In MCNP eigenvalue calculation,  $\sigma_{\text{stat}}$  is usually underestimated.
- \* An independent estimation of  $\sigma_{stat}$  is therefore necessary for fast TMC,
- \* From the 508 short runs, the first 389 were repeated with fixed nuclear data,
- ✤ 274 were then accepted due to source convergence.
- \* 9 % difference for  $k_{eff}$
- \* for generated power (f7): ratio is  $1.019 \pm 0.040$ .

Therefore, for local power, the MCNP estimation of  $\sigma_{stat}$  is good enough.



#### **Fast TMC & GRS methods on a full core:** k<sub>eff</sub> **uncertainty**



#### Fast TMC & GRS methods on a full core: generated local power



#### Fast TMC & GRS methods on a full core: generated local power



# Conclusions

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## fast TMC and GRS methods: If we can do a calculation once, we can also get nuclear data uncertainties in twice the time (or less).

